

## HIGHER PRESSURE DERIVATIVE OF BULK MODULUS AT DIFFERENT COMPRESSIONS

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### ABSTRACT

*Equation of state (EOS) Plays an important role in understanding the properties of matter under extreme conditions of pressure and temperature. In the present paper we shall consider four EOSs such as Vinet Rydberg EOS, Shanker EOS, Shanker – Kushwah EOS and Kean's EOS to study the behavior of five solids viz Ne, Cu, Al, LiH and MgO upto very high Compressions.*

**Keywords:** Equations of State, Bulk Modulus, Monoatomic and Diatomic Solids.

### Introduction

Equation of state (EOS) normally means a relationship between pressure P, Volume V and temperature T for a material [1]. An EOS can be classified into two categories, viz (i) Inverted type EOS, and (ii) non-inverted type EOS [2]. Shanker and Kushwah [3] have formulated an equation of state which can be expressed inverted as well as non-inverted forms with the help of this EOS, we can determine V/V<sub>0</sub> as a function of pressure and vice-versa. This EOS has been found most suitable for the materials for which the pressure derivative of isothermal bulk modulus at P=0 ie B<sub>0</sub> is less than four. A common feature of the phenomenological equations is that they represent the relationship between pressure and volume which can be expressed analytically involving two quantities only viz B<sub>0</sub>, B'<sub>0</sub> which are respectively the isothermal bulk modulus and its first pressure derivative both at zero pressure. Shanker et al. have developed a method within the framework of the theory of lattice Potential for ionic crystals to obtain analytical forms of EOS containing B<sub>0</sub> and B'<sub>0</sub> [4, 5]. Morse [6] presented an early quantum mechanical theory with two exponential terms to explain molecular vibration spectra. The Morse potential was first used in a finite strain theory by Slater [7]. Rydberg [8] tried the Morse potential in a study of the vibrational spectrum of molecular hydrogen. This is the potential favoured more recently and advocated as 'universal' by Vinet et al. [9]. The first true B-primed equation was developed by Keane [10]. Keane recognized the significance of the extrapolation of B to infinite pressure that is B<sub>∞</sub>.

### Formulation and Method of Analysis

We consider the following four equations:

- Vinet- Rydberg EOS
- Shanker EOS
- Shanker – Kushwah EOS
- Kean's EOS

### Vinet-Rydberg EOS

$$B_T = B_0(x)^{-\frac{2}{3}} \left[ 1 + \left\{ \eta(x)^{\frac{1}{3}} + 1 \right\} \left\{ 1 - (x)^{\frac{2}{3}} \right\} \right] \exp \left[ \eta \left\{ 1 - (x)^{\frac{1}{3}} \right\} \right]$$

$$B'_T = \frac{1}{3} \left[ \frac{(x)^{\frac{1}{3}}(1 - \eta) + 2\eta(x)^{\frac{2}{3}}}{1 + \left[ \eta(x)^{\frac{1}{3}} + 1 \right] \left[ 1 - (x)^{\frac{1}{3}} \right]} + \eta(x)^{\frac{1}{3}} + 2 \right]$$

Where  $x = \frac{V}{V_0}$  and  $\eta = \frac{3}{2}(B'_0 - 1)$

$$B_T B''_T = \frac{1}{9} (x)^{\frac{1}{3}} \left[ \frac{\left\{ (1-\eta)x^{\frac{1}{3}} + 2\eta x^{\frac{2}{3}} \right\} \left\{ (\eta-1) - 2\eta(x)^{\frac{1}{3}} \right\}}{\left\{ 1 + \left[ \eta(x)^{\frac{1}{3}} + 1 \right] \left[ 1 - (x)^{\frac{1}{3}} \right] \right\}^2} \left\{ (\eta-1) - 4\eta(x)^{\frac{1}{3}} \right\} \right] \eta \quad \dots\dots\dots(1)$$

At  $\frac{V}{V_0} = 0, B_T B''_T = 0$

**Shanker EOS**

$$B_T = B_0 x^{\frac{-1}{3}} (1 + y + y^2) \exp(ty) + \frac{4}{3} P$$

$$B'_T = \frac{4}{3} + \left( 1 - \frac{4P}{3B_T} \right) \left[ \frac{1}{3} + x \left\{ t + \frac{(1+2y)}{(1+y+y^2)} \right\} \right]$$

Where  $x = \frac{V}{V_0}, t = B'_0 - \frac{8}{3}$  and  $y = 1 - \frac{V}{V_0}$

$$B''_T = \left( 1 - \frac{4P}{3B_T} \right) \left[ \left( \frac{-x}{B} \right) \left\{ t + \frac{1+2y}{1+y+y^2} \right\} + x \left\{ \frac{2 \left( \frac{x}{B} \right)}{1+y+y^2} - \frac{(1+2y) \left\{ \frac{x}{B} + 2y \left( \frac{x}{B} \right) \right\}}{(1+y+y^2)^2} \right\} \right] - \frac{4}{3} \left( \frac{1}{B_T} - \frac{PB'_T}{B_T^2} \right) \left\{ \frac{1}{3} + x \left( t + \frac{1+2y}{1+y+y^2} \right) \right\}$$

$$B_T B''_T = \left( \frac{4P}{3B_T} - 1 \right) x \left[ \left( t + \frac{1+2y}{1+y+y^2} \right) - x \left\{ \frac{1-2y-2y^2}{(1+y+y^2)^2} \right\} \right] - \frac{4}{3} \left( 1 - \frac{PB'_T}{B_T} \right) \left\{ \frac{1}{3} + x \left( t + \frac{1+2y}{1+y+y^2} \right) \right\} \quad \dots\dots\dots(2)$$

At  $x = 0, B'_T = \frac{B_T}{P} = \frac{4}{3}$  and therefore  $B_T B''_T = 0$

**Shanker-Kushwah EOS**

$$B'_T = \frac{4(a_1 + a_2) - (a_1 + 2a_2) \left( \frac{V}{V_0} \right)}{a_1 \left[ 2 - \left( \frac{V}{V_0} \right) \right] + 2a_2 \left[ 1 - \left( \frac{V}{V_0} \right) \right]}$$

Where  $a_1 = B_0$  and  $a_2 = B_0 (B'_0 - 3)/2$

$B_0$  and  $B'_0$  are the values of  $B_T$  and  $\frac{dB_T}{dP}$  at  $V = V_0$

$$B''_T = \frac{B_T \left(\frac{V}{V_0}\right)^2 \left[-(a_1 + \lambda a_2) \frac{1}{V_0} \frac{dV}{dP}\right] - \left[B'_T \left(\frac{V}{V_0}\right)^2 + \lambda B_T \frac{V}{V_0^2} \frac{dV}{dP}\right] \left[4(a_1 + a_2) - (a_1 + \lambda a_2) \frac{V}{V_0}\right]}{B_T^2 \left(\frac{V}{V_0}\right)^4}$$

$$B_T B''_T = \frac{\left[(a_1 + 2a_2) \frac{V}{V_0}\right] - (B'_T - 2) \left[4(a_1 + a_2) - (a_1 + 2a_2) \frac{V}{V_0}\right]}{a_1 \left(2 - \frac{V}{V_0}\right) + 2a_2 \left(1 - \frac{V}{V_0}\right)} \dots \dots \dots (3)$$

At  $\frac{V}{V_0} = 0$ ,  $B'_T = B'_\infty = 2$  and therefore  $B_T B''_T = 0$

**Keane's EOS**

$$B_T = B_0 \left[1 + \frac{B'_0}{B'_\infty} (x^{-B'_\infty} - 1)\right]$$

Where  $x = \frac{V}{V_0}$

$$B'_T = B'_\infty + (B'_0 - B'_\infty) \frac{B_0}{B_T^2}$$

$$B''_T = -B'_T (B'_0 - B'_\infty) \frac{B_0}{B_T^2}$$

$$B_T B''_T = -B'_T (B'_0 - B'_\infty) \frac{B_0}{B_T} \dots \dots \dots (4)$$

**Results and Discussion**

We have calculated the values of  $B_T B''_T$  using the Vinet- Rydberg EOS, the Shanker EOS, the Shanker-Kushwah EOS and the Kean's EOS as function of  $\frac{V}{V_0}$  down to 0.1 for the solids under study taking input parameters are presented in Table 1 based on first principal reported by Hama and Suito [11] and by Anderson [12] based on experimental data. It should be mentioned that  $B'_0$  is a quantity of central importance and its value remains between three and eight for most of the solids. We have selected five solids Ne, Cu, Al LiH and MgO. For which the value of  $B'_0$  are 7.61, 5.93, 4.85, 3.51 and 4.37 respectively. The results for  $B_T B''_T$  calculated from different EOS (eqn 1 to 4) are presented in Table 2,3,4,5, and 6. The values of  $B_T B''_T$  obtained from Kean's EOS are decreasing most rapidly as compared to other EOS. Value of  $B_T B''_T$  are negative because  $B_T$  decreases with the increase in pressure and attains a constant value  $B'_T = B'_\infty$  in the limit of infinite pressure. Therefore  $B_T B''_T$  become zero in the limit  $P \rightarrow \infty$  or  $V \rightarrow 0$ . At infinity large compressions  $V \rightarrow 0, P \rightarrow \infty B'_T \rightarrow B'_\infty$  and  $\left(\frac{B_T}{P}\right) \rightarrow B'_\infty$ . Using these limits and value of  $B'_\infty$  for different EOS it can be verified that  $B_T B''_T$  vanish at  $\frac{V}{V_0} \rightarrow 0$  for all the EOS under study.

Value of input parameters for the solids based on first principles reported by Hama and Suito [11] and for NaCl reported by Anderson [12] based on experimental ultrasonic data value of  $B'_\infty$  are based on the Stacey formulation.

**Table 1**

Material	$B_0$ (GPa)	$B'_0$	$B''_0$ (GPa <sup>-1</sup> )	$B'_\infty$
Ne	6.36	7.61	2.86	5.22
Al	72.6	4.85	0.104	3.29
Cu	135	5.93	0.083	4.04
LiH	39.1	3.51	0.106	2.33
MgO	157	4.37	0.040	2.93

**Table 2**

Values of  $-B_T B_T''$  with the variation of  $V/V_0$  for Ne calculated from (a) Vinet-Rydberg EOS (b) Shanker EOS (c) Shanker-Kushwah EOS (d) Keane's EOS

$V/V_0$	(a)	(b)	(c)	(d)
1.0	17.8	13.3	37.1	18.2
0.9	9.19	8.34	13.7	7.36
0.8	5.56	6.01	6.58	3.28
0.7	3.72	4.71	3.60	1.47
0.6	2.65	3.77	2.11	0.621
0.5	2.02	2.99	1.28	0.233
0.4	1.53	2.30	0.776	0.072
0.3	1.19	1.16	0.455	0.016
0.2	0.906	0.975	0.236	0.002
0.1	0.634	0.391	0.096	0.00005

**Table 3**

Values of  $-B_T B_T''$  with the variation of  $V/V_0$  for Al calculated from (a) Vinet-Rydberg EOS (b) Shanker EOS (c) Shanker-Kushwah EOS (d) Keane's EOS

$V/V_0$	(a)	(b)	(c)	(d)
1.0	7.78	6.87	10.9	7.57
0.9	4.95	4.96	5.98	4.13
0.8	3.37	3.75	3.55	2.34
0.7	2.42	2.90	2.23	1.33
0.6	1.79	2.24	1.44	0.736
0.5	1.36	1.75	0.931	0.382
0.4	1.03	1.23	0.596	0.177
0.3	0.798	0.828	0.366	0.067
0.2	0.595	0.469	0.202	0.018
0.1	0.402	0.181	0.084	0.002

**Table 4**

Values of  $-B_T B_T''$  with the variation of  $V/V_0$  for Cu calculated from (a) Vinet-Rydberg EOS (b) Shanker EOS (c) Shanker-Kushwah EOS (d) Keane's EOS

$V/V_0$	(a)	(b)	(c)	(d)
1.0	10.4	9.39	19.4	11.2
0.9	6.58	6.36	8.99	5.42
0.8	4.26	4.69	4.86	2.78
0.7	2.95	3.62	2.85	1.44
0.6	2.15	2.84	1.75	0.716
0.5	1.62	2.20	1.10	0.328
0.4	1.24	1.64	0.687	0.130
0.3	0.956	1.23	0.412	0.040
0.2	0.720	0.569	0.224	0.008
0.1	0.494	0.255	0.093	0.0005

**Table 5**

Values of  $-B_T B_T''$  with the variation of  $V/V_0$  for LiH calculated from (a) Vinet-Rydberg EOS (b) Shanker EOS (c) Shanker-Kushwah EOS (d) Keane's EOS

$V/V_0$	(a)	(b)	(c)	(d)
1.0	4.31	3.75	3.79	4.14
0.9	3.06	3.06	2.57	2.63
0.8	2.24	2.24	1.79	1.69
0.7	1.69	1.92	1.26	1.09
0.6	1.29	1.45	0.884	0.687
0.5	0.998	1.06	0.615	0.416
0.4	0.768	0.730	0.417	0.234
0.3	0.579	0.452	0.269	0.115
0.2	0.432	0.235	0.155	0.044
0.1	0.285	0.079	0.068	0.009

**Table 6**

Values of  $-B_T B_T'$  with the variation of  $V/V_0$  for MgO calculated from (a) Vinet-Rydberg EOS (b) Shanker EOS (c) Shanker-Kushwah EOS (d) Keane's EOS

$V/V_0$	(a)	(b)	(c)	(d)
1.0	6.43	5.75	7.99	6.29
0.9	4.25	4.13	4.70	3.62
0.8	2.97	3.30	2.94	2.14
0.7	2.16	2.57	1.91	1.27
0.6	1.66	1.96	1.26	0.739
0.5	1.23	1.68	0.836	0.406
0.4	0.948	1.06	0.544	0.202
0.3	0.725	0.697	0.338	0.085
0.2	0.538	0.394	0.189	0.026
0.1	0.360	0.247	0.080	0.003

### Conclusions

In the present study we have thus found that all values of  $B_T B_T''$  are negative and tend to zero as  $\frac{V}{V_0} = 0$ .

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